

Rényi divergence and the central limit theorem

Sergey Bobkov (University of Minnesota)

based on joint work with Gennadiy Chistyakov
and Friedrich Götze (Bielefeld University)

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Rényi divergence

X, Z random elements in a measure space (Ω, μ)

P, Q distributions with densities $p = \frac{dP}{d\mu}$, $q = \frac{dQ}{d\mu}$

$0 < \alpha < \infty$

Definition (Rényi 1961, Tsallis 1998)

The Rényi divergence of P from Q of index α is

$$D_\alpha(X||Z) = D_\alpha(P||Q) = \frac{1}{\alpha - 1} \log \int \left(\frac{p}{q}\right)^\alpha q d\mu.$$

The Rényi divergence power or relative Tsallis entropy

$$T_\alpha(X||Z) = T_\alpha(P||Q) = \frac{1}{\alpha - 1} \left[\int \left(\frac{p}{q}\right)^\alpha q d\mu - 1 \right].$$

Monotone transformations:

$$D_\alpha = \frac{1}{\alpha - 1} \log (1 + (\alpha - 1) T_\alpha),$$
$$T_\alpha = \frac{1}{\alpha - 1} [e^{(\alpha-1)D_\alpha} - 1].$$

Equivalence: $D_\alpha \sim T_\alpha$ (when small)

Properties

- Independence of the dominating measure μ
- Separation: $D_\alpha(P||Q) \geq 0$, and $D_\alpha(P||Q) = 0$ if and only if $P = Q$.
- Monotonicity: The functions $\alpha \rightarrow D_\alpha(P||Q)$ and $\alpha \rightarrow T_\alpha(P||Q)$ are non-decreasing.
- Contractivity under mappings:

$$D_\alpha(S(X)||S(Z)) \leq D_\alpha(X||Z) \quad (\alpha \geq 1).$$

- Range $0 < \alpha < 1$: All D_α are comparable to each other and are metrically equivalent to the total variation $\|P - Q\|_{\text{TV}}$. Gilardoni's inequality (2010):

$$D_\alpha(P||Q) \geq \frac{\alpha}{2} \|P - Q\|_{\text{TV}}^2.$$

This extends the Pinsker inequality for the Kullback-Leibler distance ($\alpha = 1$).

- Range $\alpha \geq 1$: $D_\alpha(P||Q) < \infty \Rightarrow P \ll Q$.

Particular cases

- $\alpha = 1/2$ (Hellinger distance)
- $\alpha = 1$ (Kullback-Leibler distance, relative entropy):

$$D(X||Z) = D(P||Q) = \int p \log \frac{p}{q} d\mu.$$

- $\alpha = 2$ (χ^2 -distance, quadratic Renyi divergence):

$$D_2(X||Z) = \log \int \frac{p^2}{q} d\mu,$$
$$\chi^2(X, Z) = T_2(X||Z) = \int \frac{(p - q)^2}{q} d\mu.$$

In all cases

$$\frac{1}{2} \|P - Q\|_{\text{TV}}^2 \leq D(X||Z)$$
$$\leq D_2(X||Z) \leq \chi^2(X, Z).$$

Goodness of fit test (Karl Pearson 1900): If Q is unknown distribution with k atoms, $P = P_n$ empirical, then

$$n\chi^2(P_n, Q) \Rightarrow \chi_{k-1}^2 = \mathcal{L}(Z_1^2 + \cdots + Z_{k-1}^2)$$

Rényi divergence from Gaussian

$\Omega = \mathbb{R}^d$ with Lebesgue measure $d\mu(x) = dx$

X random vector with density $p(x)$

$Z \sim N(0, I)$ standard normal random vector with density

$$\varphi(x) = \frac{1}{(2\pi)^{d/2}} e^{-|x|^2/2}, \quad x \in \mathbb{R}^d.$$

Rényi divergence–Tsallis distance of index α are given by

$$D_\alpha(X||Z) = \frac{1}{\alpha - 1} \log \int \frac{p^\alpha}{\varphi^{\alpha-1}} dx,$$
$$T_\alpha(X||Z) = \frac{1}{\alpha - 1} \int \frac{p^\alpha}{\varphi^{\alpha-1}} dx - 1.$$

Relative entropy ($\alpha = 1$), $\mathbb{E}X = 0$, $\text{cov}(X) = I$

$$D(X||Z) = h(Z) - h(X)$$

in terms of Shannon entropy $h(X) = - \int p \log p dx$.

Pearson ($\alpha = 2$)

$$\chi^2(X, Z) = \int \frac{(p - \varphi)^2}{\varphi} dx.$$

Exponential integrability

Let $d = 1$, $Z \sim N(0, 1)$, $\beta = \frac{\alpha}{\alpha-1}$.

Note: If $D(X||Z) < \infty$, then $\mathbb{E}X^2 < \infty$.

Proposition 1. If $T_\alpha = T_\alpha(X||Z) < \infty$ for $\alpha > 1$, then X has an absolutely continuous distribution and finite moments of any order. Moreover,

$$\mathbb{E} e^{cX^2} < \infty \quad \text{for all } c < 1/(2\beta).$$

For all $t \in \mathbb{R}$,

$$\mathbb{E} e^{tX} \leq C e^{\beta t^2/2} \quad \text{with } C = (1 + (\alpha - 1) T_\alpha)^{1/\alpha}.$$

It is possible that $T_\alpha < \infty$, while $\mathbb{E} e^{\frac{1}{2\beta} X^2} = \infty$.

If p is density of X ,

$$\begin{aligned} \mathbb{E} e^{tX} &= \int_{-\infty}^{\infty} p(x) e^{tx} dx \\ &= \int_{-\infty}^{\infty} \frac{p(x)}{\varphi(x)^{1/\beta}} \cdot e^{tx} \varphi(x)^{1/\beta} dx \\ &\leq C \left(\int_{-\infty}^{\infty} e^{\beta tx} \varphi(x) dx \right)^{1/\beta} = C e^{\beta t^2/2}. \end{aligned}$$

Improved integrability for convolutions

Case $\alpha = 2$: If $\chi^2 = \chi^2(X, Z) < \infty$, then

$$\mathbb{E} e^{cX^2} < \infty \quad \text{for all } c < 1/4.$$

For all $t \in \mathbb{R}$,

$$\mathbb{E} e^{tX} \leq C e^{t^2/4} \quad \text{with } C = (1 + \chi^2)^{1/2}.$$

Proposition 2. If X_1, X_2 are independent copies of X ,

$$\mathbb{E} e^{\frac{1}{4} \left(\frac{X_1 + X_2}{\sqrt{2}} \right)^2} \leq 2(1 + \chi^2).$$

For general $\alpha > 1$ with conjugate $\beta = \frac{\alpha}{\alpha-1}$, we need $k \geq \alpha$ normalized convolutions to include the critical coefficient $c = 1/(2\beta)$: If $T_\alpha = T_\alpha(X||Z)$ is finite, then

$$\mathbb{E} e^{\frac{1}{2\beta} Z_k^2} \leq 2^k (1 + (\alpha - 1) T_\alpha)^{\frac{k}{\alpha}}.$$

Moreover, if $T_\alpha(Z_k||Z) \rightarrow 0$, then

$$\mathbb{E} e^{\frac{1}{2\beta} Z_k^2} \rightarrow \mathbb{E} e^{\frac{1}{2\beta} Z^2} = \sqrt{\pi(\alpha - 1)}.$$

Proof. Use of Plancherel theorem in case $\alpha = 2$ (Weierstrass transform for $\alpha > 1$).

Exponential series and normal moments

Question: How to connect $\chi^2(X, Z)$ to the moments of X ?
Exponential orthogonal series (Cramér):

$$p(x) = \varphi(x) \sum_{k=0}^{\infty} \frac{c_k}{k!} H_k(x)$$

converges in $L^2(\mathbb{R}, \frac{dx}{\varphi(x)})$ if and only if $\sum_{k=0}^{\infty} c_k^2 < \infty$.

Fourier coefficients (normal moments of X):

$$c_k = \int_{-\infty}^{\infty} H_k(x) p(x) dx = \mathbb{E} H_k(X) = \mathbb{E} (X + iZ)^k.$$

In particular, $c_0 = 1$, $c_1 = \mathbb{E}X$, $c_2 = \mathbb{E}X^2 - 1$.

Taylor series around zero for the characteristic function:

$$f(t) = \mathbb{E} e^{itX} = e^{-t^2/2} \sum_{k=0}^n \frac{c_k}{k!} (it)^k + o(|t|^n).$$

Proposition 3. If $\chi^2(X, Z) < \infty$, then

$$\chi^2(X, Z) = \sum_{k=1}^{\infty} \frac{1}{k!} (\mathbb{E} H_k(X))^2.$$

Conversely, if X has all moments and the series is convergent, then $\chi^2(X, Z) < \infty$. In particular, X has density.

CLT for strong metrics

X, X_1, X_2, \dots i.i.d. random variables, $\mathbb{E}X = 0$, $\mathbb{E}X^2 = 1$,

$$Z_n = \frac{X_1 + \dots + X_n}{\sqrt{n}} \quad (n = 1, 2, \dots)$$

CLT: as $n \rightarrow \infty$

$$F_n(x) = \mathbb{P}\{Z_n \leq x\} \rightarrow \Phi(x) = \int_{-\infty}^x \varphi(y) dy.$$

Total variation distance. Prokhorov (1952):

$$\|F_n - \Phi\|_{\text{TV}} \rightarrow 0 \iff \|F_{n_0} - \Phi\|_{\text{TV}} < 2,$$

F_n has an absolutely continuous component for some $n = n_0$ (in particular, if X has density).

Kullback-Leibler distance (relative entropy). Barron (1986):

$$D(Z_n || Z) \rightarrow 0 \iff D(Z_{n_0} || Z) < \infty$$

for some $n = n_0$. In particular, when X has density p such that $\int_{-\infty}^{\infty} p(x) \log p(x) dx < \infty$.

Rates, Berry-Esseen bounds, the non-i.i.d. case: Linnik (1959), Sirazhdinov, Mamatov (1962), Artstein, Ball, Barthe, Naor (2004), Barron, Johnson (2004), B-C-G (2013-2016), Toscani (2016), Bally, Caramellino (2016).

CLT for χ^2 distance

Fomin (1982): Suppose that X has a compactly supported, symmetric, piecewise differentiable density p such that the coefficients in

$$p(x) = \varphi(x) \sum_{k=1}^{\infty} \frac{\sigma_k}{2^k k!} H_{2k}(x)$$

satisfy $\sup_{k \geq 2} \sigma_k < 1$. Then $\chi^2(Z_n, Z) = O(\frac{1}{n})$ as $n \rightarrow \infty$.

Example: Uniform distribution on $(-\sqrt{3}, \sqrt{3})$.

Theorem 1. $\chi^2(Z_n, Z) \rightarrow 0$, if and only if $\chi^2(Z_n, Z) < \infty$ for some $n = n_0$, and

$$\mathbb{E} e^{tX} < e^{t^2} \quad \text{for all } t \neq 0.$$

Remark. It is possible that $\mathbb{E} e^{tX} < e^{t^2}$ for all $t \neq 0$ except for one $t_0 > 0$. Consider $X = a\xi + bZ$ assuming that ξ takes values q and $-p$ with probabilities p and q such that

$$\frac{p - q}{\log p - \log q} > pq.$$

Edgeworth-type expansion

If $\chi^2(Z_n, Z) \rightarrow 0$, then as $n \rightarrow \infty$

$$\chi^2(Z_n, Z) = \sum_{j=1}^{s-2} \frac{c_j}{n^j} + O\left(\frac{1}{n^{s-1}}\right)$$

for every fixed $s = 3, 4, \dots$ with c_j certain polynomials in the moments $\alpha_k = \mathbb{E}X^k$, $k = 3, \dots, j + 2$.

Case $s = 3$:

$$\chi^2(Z_n, Z) = \frac{\alpha_3^2}{6n} + O\left(\frac{1}{n^2}\right),$$

Case $\alpha_3 = 0$, $s = 4$:

$$\chi^2(Z_n, Z) = \frac{(\alpha_4 - 3)^2}{24n^2} + O\left(\frac{1}{n^3}\right).$$

CLT for Renyi divergence

X, X_1, X_2, \dots i.i.d. random vectors in \mathbb{R}^d , with $\mathbb{E}X = 0$ and identity covariance.

Denote by $\alpha^* = \frac{\alpha}{\alpha-1}$ the conjugate index for $\alpha > 1$, and by Z a standard normal random vector in \mathbb{R}^d .

Theorem 2. $D_\alpha(Z_n||Z) \rightarrow 0$ if and only if $D_\alpha(Z_n||Z) < \infty$ for some $n = n_0$, and

$$\mathbb{E} e^{\langle t, X \rangle} < e^{\alpha^* |t|^2 / 2} \quad \text{for all } t \in \mathbb{R}^d, t \neq 0.$$

In this case,

$$D_\alpha(Z_n||Z) = O(1/n),$$

and even $D_\alpha(Z_n||Z) = O(1/n^2)$, if the distribution of X is symmetric.

In fact, all distances have similar rates

$$D_\alpha(Z_n||Z) \sim T_\alpha(Z_n||Z) \sim \frac{\alpha}{2} \chi^2(Z_n, Z)$$

once they tend to zero.

Examples

- Uniform distribution.

Let $X \sim (-\sqrt{3}, \sqrt{3})$. It has Laplace transform

$$\mathbb{E} e^{tX} = \frac{\sinh(t\sqrt{3})}{t\sqrt{3}} < e^{t^2/2}, \quad t \in \mathbb{R} \quad (t \neq 0),$$

with first moments $\alpha_2 = 1$, $\alpha_3 = 0$, $\alpha_4 = \frac{9}{5}$. Therefore,

$$\begin{aligned} \chi^2(Z_n, Z) &= \frac{3}{50n^2} + O\left(\frac{1}{n^3}\right), \\ D_\alpha(Z_n || Z) &= \frac{\alpha}{2} \chi^2(Z_n, Z) + O\left(\frac{1}{n^3}\right). \end{aligned}$$

- Log-concave probability distributions on \mathbb{R}^d .

Let X have density $p(x) = e^{-V(x)}$ with mean zero, identity covariance and such that $V''(x) \geq cI$ for some $c > 0$ (Bakry-Emery criterion, necessarily $c \leq 1$). Then

$$\mathbb{E} e^{tg(X)} \leq e^{t^2/(2c)}, \quad t \in \mathbb{R}.$$

for any g on \mathbb{R}^d such that $\|g\|_{\text{Lip}} \leq 1$ and $\mathbb{E} g(X) = 0$. Hence

$$D_\alpha(Z_n || Z) \rightarrow 0 \text{ as } n \rightarrow \infty, \text{ whenever } \alpha < \frac{1}{1-c}.$$

Necessity part in Theorem 2 (preparation)

The characteristic function $f(t) = \mathbb{E} e^{itX}$ is entire on \mathbb{C} , and $f(iy) = \mathbb{E} e^{-yX}$.

Lemma 1. If $\lim D_\alpha(Z_n||Z) = 0$, then for all $y \in \mathbb{R}$

$$f(iy) \leq e^{\beta y^2/2}$$

and for any integer $k \geq \alpha/2$,

$$\lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} f(iy/\sqrt{kn})^{2kn} e^{-\beta y^2} dy = \sqrt{\pi(\alpha - 1)}.$$

Proof. By Proposition 1 applied to Z_n in place of X ,

$$f(iy/\sqrt{n})^n = \mathbb{E} e^{-yZ_n} \leq C_n e^{\beta y^2/2}$$

with $C_n = (1 + (\alpha - 1) T_\alpha(Z_n||Z))^{1/\alpha}$. After change of variable

$$f(iy) \leq C_n^{1/n} e^{\beta y^2/2} \rightarrow e^{\beta y^2/2}.$$

Since $f(\frac{t}{\sqrt{n}})^n$ is the ch.f. of Z_n , the integral in Lemma is

$$\begin{aligned} \int_{-\infty}^{\infty} (\mathbb{E} e^{-yZ_{nk}})^2 e^{-\beta y^2} dy &= \int_{-\infty}^{\infty} \mathbb{E} e^{-\sqrt{2}y Z_{2nk}} e^{-\beta y^2} dy \\ &= \sqrt{\frac{\pi}{\beta}} \mathbb{E} e^{\frac{1}{2\beta} Z_{2nk}^2} \rightarrow \sqrt{\frac{\pi}{\beta}} \mathbb{E} e^{\frac{1}{2\beta} Z^2}, \end{aligned}$$

by Proposition 2 on last step.

Necessity part in Theorem 2

If $D_\alpha(Z_n||Z) \rightarrow 0$, then, by Lemma 1,

$$\psi(y) = f(iy) e^{-\beta y^2/2} \leq 1 \quad \text{for all } y \in \mathbb{R}.$$

Need to show:

$$\psi(y) < 1 \quad \text{for all } y \neq 0.$$

Fix $k \geq \alpha/2$ and $\delta > 0$ (small), and decompose

$$\begin{aligned} \int_{-\infty}^{\infty} f(iy/\sqrt{nk})^{2nk} e^{-\beta y^2} dy &= I_1 + I_2 \\ &= \left(\int_{|y| \leq \delta\sqrt{nk}} + \int_{|y| > \delta\sqrt{nk}} \right) f(iy/\sqrt{nk})^{2nk} e^{-\beta y^2} dy. \end{aligned} \quad (1)$$

Write

$$g(t) = \log f(t) = -\frac{1}{2}t^2 + \sum_{m=3}^{\infty} a_m t^m.$$

Since $\sum_{m=3}^{\infty} |a_m t^m| \leq c|t|^3$ for $|t| \leq r$, $r > 0$ small, so,

$$f(iy/\sqrt{nk})^{2nk} = \exp\{y^2 + \theta y^3/\sqrt{n}\}, \quad y \in [-r\sqrt{nk}, r\sqrt{nk}],$$

where $|\theta| \leq c$. Assuming $\delta \leq \min\{r, (\beta - 1)/(2c\sqrt{k})\}$,

$$I_1 = \int_{|y| \leq \delta\sqrt{nk}} e^{-(\beta-1)y^2 + \theta y^3/\sqrt{n}} dy.$$

Here $\theta y^3 / \sqrt{n}$ may be removed at the expense of $O(\frac{1}{\sqrt{n}})$. Hence

$$I_1 = \int_{|y| \leq \delta \sqrt{nk}} e^{-(\beta-1)y^2} dy + O\left(\frac{1}{\sqrt{n}}\right) = \sqrt{\pi(\alpha-1)} + O\left(\frac{1}{\sqrt{n}}\right).$$

Applying this in (1), we have $I_2 \rightarrow 0$, or equivalently

$$\int_{|u| > \delta} \psi(u)^{2nk} du = \int_{|u| > \delta} (f(iu) e^{-\beta u^2/2})^{2nk} du = o\left(\frac{1}{\sqrt{n}}\right)$$

for any sufficiently small $\delta > 0$ and hence for any $\delta > 0$.

Assume $\psi(u_0) = 1$ for some $u_0 > 0$, which implies $\psi'(u_0) = 0$.

Hence the power series representation at this point

$$\psi(u) - 1 = c_l(u - u_0)^l + \sum_{j=l+1}^{\infty} c_j(u - u_0)^j$$

starts with $c_l \neq 0$ for some $l \geq 2$. Since $\psi(u) \leq 1$ for all $u \in \mathbb{R}$, necessarily $l = 2m$ and $c_l < 0$. Hence,

$$\psi(u) \geq 1 - b_1(u - u_0)^{2m} \geq e^{-b_0(u - u_0)^{2m}}.$$

for $|u - u_0| \leq r_0 < u_0$ with some constants $b_1, b_0 > 0$. Choosing $\delta = u_0 - r_0$, this neighborhood is contained in (δ, ∞) , and

$$\begin{aligned} \int_{|u| > \delta} \psi(u)^{2nk} du &\geq \int_{|u - u_0| < \delta} \exp\{-2nk \cdot c_0(u - u_0)^{2m}\} du \\ &= 2 \int_0^\delta \exp\{-2nk \cdot c_0 x^{2m}\} dx \geq \frac{c}{n^{1/(2m)}}. \end{aligned}$$

Pointwise upper bounds on densities

The following observation holds without assuming that X has mean zero and variance one. Let f be the characteristic function of X , and define

$$\psi(u) = f(iu) e^{-\beta u^2/2} = \mathbb{E} e^{-uX} e^{-\beta u^2/2}, \quad u \in \mathbb{R},$$

where $\beta = \frac{\alpha}{\alpha-1}$.

Question: If $T_\alpha(X||Z) < \infty$, can we bound the density $p(x)$ of X pointwise? **Answer:** No. However, assume $T_\alpha(Z_{n_0}||Z) < \infty$.

Lemma 2. For any $n \geq n_\beta = \max(\beta, 2) n_0$, the normalized sum Z_n has a continuous bounded density p_n satisfying

$$p_n(x) \leq \frac{A_\alpha \sqrt{n}}{\sqrt{2\pi n_0}} e^{-x^2/(2\beta)} \psi\left(-\frac{x}{\beta\sqrt{n}}\right)^{n-n_\beta}, \quad x \in \mathbb{R}.$$

In particular, there exist $x_0 > 0$ and $\delta \in (0, 1)$ such that, for all n large enough,

$$p_n(x) \leq \delta^n e^{-x^2/(2\beta)} \psi\left(-\frac{x}{\beta\sqrt{n}}\right)^{n/2}, \quad |x| \geq x_0\sqrt{n}.$$

Proof. Use contour integration.

Proof of Lemma 2

Let $\alpha = 2$, $n_0 = 1$, so that f_n integrable for $n \geq 2$ and

$$p_n(x) = e^{yx} \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-itx} f((t + iy)/\sqrt{n})^n dt$$

for any fixed $y > 0$. Let $x < 0$. Using $|f(t + iy)| \leq f(iy)$ and changing variable,

$$p_n(x) \leq e^{yx} f\left(\frac{iy}{\sqrt{n}}\right)^{n-2} \frac{\sqrt{n}}{2\pi} \int_{-\infty}^{\infty} \left|f\left(t + \frac{iy}{\sqrt{n}}\right)\right|^2 dt.$$

The function $t \rightarrow f(t + iy/\sqrt{n}) = \mathbb{E} e^{itX - yX/\sqrt{n}}$ is the Fourier transform of $e^{-yu/\sqrt{n}} p(u)$ and

$$e^{-2yu/\sqrt{n}} p(u)^2 = (e^{-2yu/\sqrt{n}} \varphi(u)) \frac{p(u)^2}{\varphi(u)} \leq \frac{1}{\sqrt{2\pi}} e^{\frac{2y^2}{n}} \frac{p(u)^2}{\varphi(u)}.$$

By Plancherel,

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \left|f\left(t + \frac{iy}{\sqrt{n}}\right)\right|^2 dt \leq \frac{1}{\sqrt{2\pi}} e^{2y^2/n} (1 + \chi^2).$$

Hence

$$\begin{aligned} p_n(x) &\leq \sqrt{\frac{n}{2\pi}} (1 + \chi^2) e^{yx+2y^2/n} f(iy/\sqrt{n})^{n-2} \\ &= \sqrt{\frac{n}{2\pi}} (1 + \chi^2) e^{yx+y^2} \psi(y/\sqrt{n})^{n-2}. \end{aligned}$$

Choose $y = -x/2$.

Sufficiency part in Theorem 2

Let X, X_1, X_2, \dots be i.i.d., $\mathbb{E}X = 0$, $\mathbb{E}X^2 = 1$, with ch.f. $f(t) = \mathbb{E} e^{itX}$. As before, put

$$\psi(u) = f(iu) e^{-\beta u^2/2}, \quad \beta = \frac{\alpha}{\alpha - 1}, \quad Z \sim N(0, 1).$$

Assuming that $\psi(u) < 1$ for all $u \neq 0$, we need to show that

$$Z_n = \frac{X_1 + \dots + X_n}{\sqrt{n}}$$

satisfy $T_\alpha(Z_n || Z) \rightarrow 0$ as long as $T_\alpha(Z_{n_0} || Z) < \infty$ for some n_0 . By Lemma 2, Z_n have densities p_n which are continuous and bounded whenever $n \geq n_\beta$.

Using Edgeworth expansions, the integrals

$$I_0 = \int_{|x| \leq M_n} \frac{p_n(x)^\alpha}{\varphi(x)^{\alpha-1}} dx \quad \text{with} \quad M_n = \sqrt{2(s-1) \log n}$$

admit an asymptotic expansion in powers of $1/n$ up to $1/n^{s-1}$. So, it remains to bound the integral of $p_n^\alpha / \varphi^{\alpha-1}$ over $|x| > M_n$ by a polynomially small quantity. In fact, for any large enough $s \geq 3$ and some constant $\kappa > 0$,

$$\int_{|x| > M_n} \frac{p_n(x)^\alpha}{\varphi(x)^{\alpha-1}} dx = O\left(\frac{1}{n^{\kappa s}}\right), \quad n \rightarrow \infty.$$

For definiteness, let $x < -M_n$, and define

$$I_1 = \int_{-\infty}^{-x_0\sqrt{n}} \frac{p_n(x)^\alpha}{\varphi(x)^{\alpha-1}} dx, \quad I_2 = \int_{-x_0\sqrt{n}}^{-x_1\sqrt{n}} \frac{p_n(x)^\alpha}{\varphi(x)^{\alpha-1}} dx,$$

$$I_3 = \int_{-x_1\sqrt{n}}^{-M_n} \frac{p_n(x)^\alpha}{\varphi(x)^{\alpha-1}} dx$$

with parameters $0 < x_1 < x_0$ and assuming that $M_n < x_1\sqrt{n}$ (otherwise, $I_3 = 0$).

By Lemma 2, for all large n , with some $\delta \in (0, 1)$, $x_0 > 0$,

$$I_1 \leq (2\pi)^{\frac{\alpha-1}{2}} \delta^{\alpha n} \int_{-\infty}^{-x_0\sqrt{n}} \psi\left(-\frac{x}{\beta\sqrt{n}}\right)^{\alpha n/2} dx$$

$$\leq (2\pi)^{\frac{\alpha-1}{2}} \delta^{\alpha n} \beta\sqrt{n} \int_{-\infty}^{\infty} \psi(u)^m du, \quad m \leq \frac{\alpha n}{2},$$

where on the last step we used $\psi \leq 1$. The last integral is convergent whenever $m = kn_0$, $k \geq \alpha$. Hence

$$I_1 \leq C\delta_1^n \quad (n \geq n_1)$$

with some constants $C > 0$, $x_0 > 0$ and $\delta < \delta_1 < 1$, depending on the density p only.

By assumption, $\delta_2 = \max_{-x_0 \leq u \leq -x_1} \psi(u) < 1$, and Lemma 2

yields

$$\begin{aligned}
I_2 &\leq A_\alpha n^{\alpha/2} \int_{-x_0\sqrt{n}}^{-x_1\sqrt{n}} \psi\left(-\frac{x}{\beta\sqrt{n}}\right)^{n-n_\beta} dx \\
&= A_\alpha^\alpha \beta n^{(\alpha+1)/2} \int_{-x_0/2}^{-x_1/2} \psi(u)^{n-n_\beta} du \\
&\leq A_\alpha^\alpha \beta n^{(\alpha+1)/2} (x_0 - x_1) \delta_2^{n-n_\beta}
\end{aligned}$$

which again decays exponentially fast like I_1 .

Near zero, $h(u) = \log f(iu) \sim \frac{1}{2} u^2$, hence $|h(u)| \leq \frac{1+\beta}{4} |u|^2$ in some disc $|u| \leq r$, when r is sufficiently small, implying $|f(iu)| \leq e^{(1+\beta)|u|^2/4}$. Hence

$$\psi(u) \leq e^{-\frac{1}{4}(\beta-1)|u|^2}, \quad |u| \leq r,$$

and

$$\begin{aligned}
\psi\left(-\frac{x}{\beta\sqrt{n}}\right)^{n-n_\beta} &\leq \psi\left(-\frac{x}{\beta\sqrt{n}}\right)^{n/2} \\
&\leq \exp\left\{-\frac{\beta-1}{4} \frac{x^2}{2\beta^2}\right\} = e^{-x^2/(8\alpha\beta)}
\end{aligned}$$

for all $n \geq 2n_\beta$ and $-\beta r\sqrt{n} < x < 0$. Therefore, By Lemma 2, in this interval

$$\frac{p_n(x)^\alpha}{\varphi(x)^{\alpha-1}} \leq A_\alpha^\alpha n^{\alpha/2} e^{-x^2/(8\beta)},$$

which results with $x_1 = \beta r$ in the bound

$$\begin{aligned} I_3 &\leq A_\alpha^\alpha n^{\alpha/2} \int_{-x_1\sqrt{n}}^{-M_n} e^{-x^2/(8\beta)} dx \\ &\leq \sqrt{2\pi\beta} A_\alpha^\alpha n^{\alpha/2} e^{-M_n^2/(8\beta)} = \sqrt{2\pi\beta} A_\alpha^\alpha n^{-\left(\frac{s-1}{4\beta} - \frac{\alpha}{2}\right)}. \end{aligned}$$

Collecting these bounds, we obtain that $I_1+I_2+I_3 = o(n^{-s/8\beta})$ for a sufficiently large s .